

Chemistry 4000/5000/7001, Fall 2012, Assignment 3 Solutions

1. (a)

$$hc\tilde{\nu}_0 = \hbar\omega_0 = \frac{h}{2\pi}\omega_0$$
$$\therefore \omega_0 = 2\pi c\tilde{\nu}_0$$

(b)

$$\omega_0 = 2\pi(2.997\,924\,58 \times 10^8 \text{ m/s})(100 \text{ cm m}^{-1})(2990.95 \text{ cm}^{-1})$$
$$= 5.633\,91 \times 10^{14} \text{ s}^{-1}$$

(c)

$$Q = \left[1 - \exp\left(-\frac{\hbar\omega_0}{k_B T}\right) \right]^{-1}$$
$$= \left[1 - \exp\left(-\frac{(1.054\,571\,73 \times 10^{-34} \text{ J s})(5.633\,91 \times 10^{14} \text{ s}^{-1})}{(1.380\,6488 \times 10^{-23} \text{ J K}^{-1})(293.15 \text{ K})}\right) \right]^{-1}$$
$$= 1.000\,000\,4$$

(I don't know how many significant figures this number really has, but I just wanted to show that it was just a bit bigger than 1.)

$$P(v=0) = \frac{1}{Q} \exp\left(-\frac{\hbar\omega_0 v}{k_B T}\right)$$
$$= \frac{e^0}{1.000\,000\,4} = 0.999\,999\,6.$$

(Again, it's too many significant figures, but the point is to illustrate that the probability that the ground state is occupied is extremely close to 1 at room temperature.)

(d) See figure 1 for my graph. At low temperatures, $Q \approx 1$, meaning that the ground state is the only one with a significant population. Values of $Q > 1$ indicate that excited states are also populated. The temperature where we think this effect is "significant" is somewhat arbitrary. Any temperature between 1000 and 2000 K would be a sensible answer.

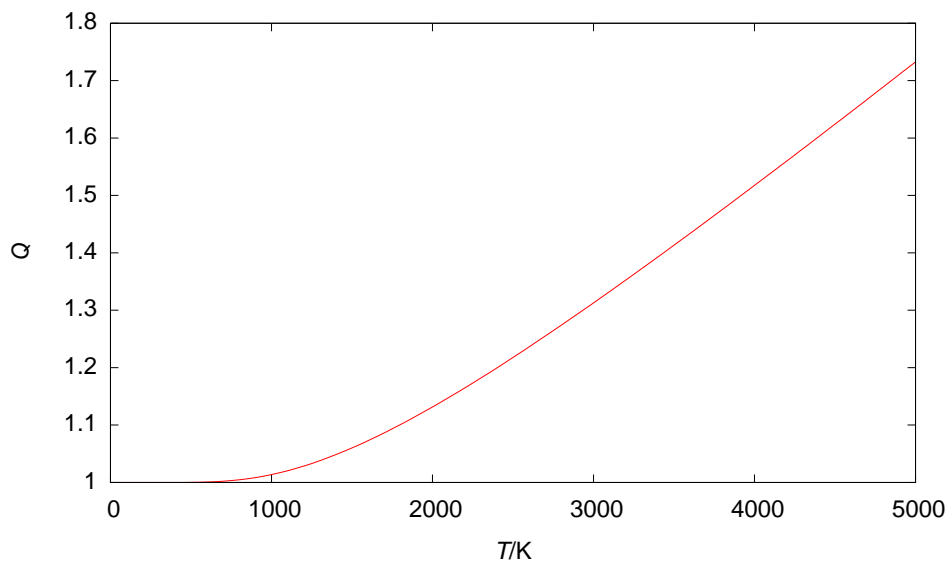


Figure 1: $Q(T)$ for the vibrational degree of freedom of HCl

2. We have the equation

$$p(\epsilon) = \frac{1}{Q} g(\epsilon) \exp\left(-\frac{\epsilon}{k_B T}\right) d\epsilon$$

for the probability density $p(\epsilon)$. This density must integrate out to 1 over the allowed energy region (\mathcal{A}), i.e.

$$\int_{\mathcal{A}} \frac{1}{Q} g(\epsilon) \exp\left(-\frac{\epsilon}{k_B T}\right) d\epsilon = 1.$$

Note that the partition function depends only on T , so we can pull it out of the integral:

$$\begin{aligned} \frac{1}{Q} \int_{\mathcal{A}} g(\epsilon) \exp\left(-\frac{\epsilon}{k_B T}\right) d\epsilon &= 1. \\ \therefore Q &= \int_{\mathcal{A}} g(\epsilon) \exp\left(-\frac{\epsilon}{k_B T}\right) d\epsilon. \end{aligned}$$

3. The value of the translational partition function answers this question. The only catch is that we have to do a whole bunch of unit conversions

in order to get our answer.

$$\begin{aligned} m &= 2(1.007\,825\,032\,07\text{ u}) = 2.015\,650\,064\,14\text{ u} \\ &\equiv \frac{2.015\,650\,064\,14\text{ g mol}^{-1}}{(1000\text{ g kg}^{-1})(6.022\,141\,29 \times 10^{23}\text{ mol}^{-1})} \\ &= 3.347\,065\,38 \times 10^{-27}\text{ kg} \end{aligned}$$

(This is one of those rare cases where the number of significant digits in the answer is limited by the number of significant digits in a fundamental constant, due in this case to the exquisite accuracy with which we are able to weigh atoms.)

$$\begin{aligned} V &= \frac{1.050\text{ L}}{1000\text{ L m}^{-3}} = 1.050 \times 10^{-3}\text{ m}^3. \\ Q_{\text{tr}} &= \frac{V}{h^3} (2\pi m k_B T)^{3/2} \\ &= (1.050 \times 10^{-3}\text{ m}^3) (6.626\,069\,57 \times 10^{-34}\text{ J s})^{-3} \\ &\quad \times [2\pi (3.347\,065\,38 \times 10^{-27}\text{ kg}) (1.380\,6488 \times 10^{-23}\text{ J K}^{-1}) (293.15\text{ K})]^{3/2} \\ &= 2.834 \times 10^{27}. \end{aligned}$$

Needless to say, this is a truly colossal number.

4. The mass of a deuterium molecule is approximately 4 u, about twice as much as the mass of an ordinary hydrogen molecule. Since Q_{tr} depends on $m^{3/2}$, the number of translational levels accessible should increase by a factor of $2^{3/2} \approx 2.828$, so the number of translational states should be approximately $(2.828)(2.834 \times 10^{27}) = 8.017 \times 10^{27}$. This happens because the particle-in-a-box energy levels depend inversely on m , i.e. the levels shift to lower energies for larger values of m .